增压式真空预压两种固结解析解的对比分析

沈宇鹏^{1,2},董淑海³,王云超¹,金亚伟⁴,李鑫宇¹

(1. 北京交通大学 轨道工程北京市重点实验室,北京 100044; 2. 北京交通大学 北京市轨道交通线路安全与防灾工程技术研究中心,北京 100044; 3. 中交第三航务工程勘察设计院有限公司,上海 200032; 4. 江苏鑫泰岩土科技有限公司,江苏 宜兴 214267)

摘要:增压式真空预压能够克服常规预压淤堵问题,但施加增压后的土体固结规律缺少理论依据,基于土体等应 变理论,结合 Newman-Carrilo 方程,利用积分变换法,提出增压式真空预压单层均质土地基固结度解析解,并与 分离变量法的固结结果进行对比分析,最后给出便于工程实践使用的图表。结果表明:随着时间因子 T_h 的增大, 固结度总体呈快速增长的趋势;随着井径比 n 增大,固结速率减缓;积分变换法与分离变量法求解得到的解析解 相差甚小,认为两种方法均可以完成增压式真空预压固结解析解的推导。

关键词: 土力学; 增压式真空预压; 积分变换法; 固结理论; 解析解

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Comparative analysis between integral-transform solution and separate variable method for consolidation of the ground by air-boosted vacuum preloading

SHEN Yupeng^{1, 2}, DONG Shuhai³, WANG Yunchao¹, JIN Yawei⁴, LI Xinyu¹

- (1. Key Laboratory of Track Engineering, Beijing Jiaotong University, Beijing 100044, China; 2. Beijing Engineering and Technology Research Center of Rail Transit Line Safety and Disaster Prevention, Beijing Jiaotong University, Beijing 100044, China;
- 3. CCCC Third Harbor Consultants Co., Ltd., Shanghai 200032, China; 4. Jiangsu Xintai Geotechnical Technology Co., Ltd., Yixin, Jiangsu 214267, China)

Abstract: Air-boosted vacuum preloading can overcome the problem of clogging of classical vertical drains, but there is lack of theoretical basis for the consolidation of air-boosted vacuum preloading method. The analytical solution of consolidation of single-layer homogeneous soil with air-boosted vacuum preloading is proposed which is based on the assumption of equal strain of soil, Newman-Carrilo equation and the intergral-transform method. And the results were compared the analytical solution with that from the separation variable method. At last, the engineering practical charts and tables are put forward to actual reference. The results show that: With the increase of T_h , the total average consolidation degree grows rapidly. With the increase of n_h , the consolidation rate decreases; the analytical solution obtained by integral-transform method and separation variable method have little difference, which can be used to complete the derivation of the analytical solution of air-boosted vacuum preloading consolidation.

Key words: soil mechanics; air-boosted vacuum preloading; intergral-transform method; consolidation theory; theoretical solution

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作者简介: 沈宇鹏(1977-), 男, 2007 年于北京交通大学道路与铁道工程专业获博士学位, 现任副教授, 主要从事路基工程方面的教学与研 究工作。E-mail: ypshen@bjtu.edu.cn

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1 引言

增压式真空预压由金亚伟等^[1]提出,该方法基于常规真空预压加固机制,而在塑料排水板之间设置增压管,相当于在排水板影响半径范围土体外施加径向压力,增加了土体与塑料排水板之间的压力差,从而加速了土体中自由水向排水板的定向流动,缩短了土体固结时间;该方法克服了常规真空预压排水板易淤堵的缺陷,已经在多个工程中得到了成功应用。目前增压式真空预压的固结理论尚不完善,该方法的固结理论研究成果较少。

有关增压式真空预压的研究工作,主要局限于该方法的工程应用。扬子江等^[2]结合珠海西站软基处理工程实例,发现了增压式真空预压具有固结速度快、效果好等优点。沈宇鹏等^[3-4]在站场地基处理中采用了增压式真空预压,验证了增压式真空预压能够有效地减小站场地基工后沉降,加速软土地基固结来缩短工期。沈宇鹏等^[5]利用有限元软件开展了增压式真空预压设计参数的优化研究,并提出了该方法的优化设计参数。J. Wang等^[6]验证了增压式真空预压技术在处理疏浚土地基方面的优越性。Y. Q. Cai 等^[7]通过将增压管改为 PVD 增压器,提出了一种加固深层海相黏土更为有效的处理方法。

目前,积分变换法主要用于成层地基的一维固 结的求解。陈根媛^[8]应用 Laplace 变换、初参数法以 及矩阵法得到了任意多层地基的一维固结解析解。 王新辉等[9]利用 Laplace 变换进行了二级加荷下的 双层地基一维固结问题。韦锋[10]将填石路基简化 为双层地基模型,借助 Laplace 变换求解了二级加 荷的双层地基一维固结问题。艾智勇和王全胜[11]结 合 Laplace 变换, 推导出 Laplace 和 Hankel 变换域 和物理域内土层任意位置的孔隙水压力、位移、应 力解析解。艾智勇等[12]通过引入中间变量和 Laplace- Fourier 变换求解出渗透各向异性的成层地 基三维 Biot 固结问题。G. J. Chen[13]利用状态矢量法 与 Laplace-Hankel 积分变换,求解了多孔弹性和多 层半空间的固结解析解。Z. Y. Ai 等[14]通过应用 Fourier 展开式,Hankel 变换和 Laplace 变换,提出 了具有各向异性渗透性和可压缩的多层多孔弹性材 料的非轴对称固结解。

积分变换法在增压式真空预压软土地基固结方 面研究较少,论文基于等应变条件,采用积分变化 法,考虑涂抹效应和井阻效应,提出增压式真空预 压单层均质土地基在任意荷载作用下单层均质土地 基的固结解析解,并与分离变量法求得解析解对比 分析。

2 单一土层的定值问题

2.1 基本假定

参照 R. A. Barron 等^[15-17]提出的研究理论和基本假定,本文采用的基本假定如下:

(1)等应变条件成立,即同一深度上任意一点的竖向变形是相等的;(2)不考虑侧向变形;(3)忽略地基土的竖向渗流,仅考虑径向渗流;(4)未扰动区与涂抹区土体除渗透系数不同外,其他性质相同;(5)忽略排水板内径向渗流,排水板内孔压沿径向不变;(6)任意深度处从土体中沿板周流入排水板的水量等于从排水板中流出的水量的增量;(7)围压 p(t)随时间变化,其加载曲线如图 1 所示,为两段式加载,t₁时刻前匀速加载至 p',保持加压值稳定至 t₂时刻;(8)将排水板等效成砂井,并按下式^[18]计算排水板等效砂井半径:

$$d_{\rm w} = \zeta \frac{2(b+d)}{\pi} \tag{1}$$

式中: d_w 为排水板直径(m); ζ 为换算系数, 取 0.75~ 1.0,一般取 1.0; b, d 分别为塑料排水板宽度和厚度(m)。

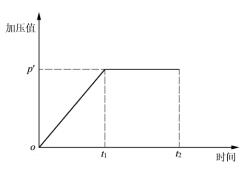


图 1 增压随时间变化曲线

Fig.1 Curve of air-boost changes over time

2.2 计算模型

增压式真空预压技术是在真空预压系统中加设一套由插入土体内部的增压管组成的径向增压系统,从而实现在软土地基中施加径向压力的效果。增压式真空预压技术可简化为图 2^[19]所示。增压式真空预压法的排水路径如图 3^[20]所示,从图中可以看出:增压式真空预压技术加固软土地基过程中,一方面是由于真空预压在排水板内部产生的负压发挥作用,另一方面则是增压管在土体内部进行径向

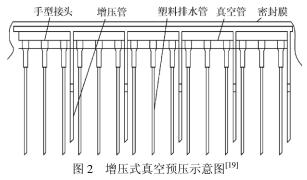


Fig.2 Schematic diagram of air-boosted vacuum preloading^[19]

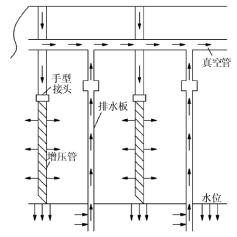


图 3 增压式真空预压法排水过程[20]

Fig.3 Air-boosted vacuum preloading drainage process^[20]

加压,加大了土体中心与排水板之间的压力差,从 而促使土体中自由水向排水板快速定向流动。

增压式真空预压排水板地基固结的理论计算模型如图 4 所示,单向排水,底面不透水。图 4 中 H 为排水板长度(m),R 为排水板影响范围(m), r_s 为涂抹区半径(m); r_w 为排水板半径(m);p(t), p_0 分别为增压管压力等效荷载和真空预压等效均布荷载(MPa); k_h , k_v 和 k_s 分别为未扰动区土体水平、竖向和涂抹区水平的渗透系数(m/s)。

2.3 等应变条件下径向固结方程

径向固结基本方程为

$$\frac{\partial \varepsilon_{v}}{\partial t} = -\frac{1}{E_{s}} \frac{\partial \overline{u_{r}}}{\partial t}$$
 (2)

$$-\frac{k_{s}}{\gamma_{w}} \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u_{r}}{\partial r} \right) = \frac{\partial \varepsilon_{v}}{\partial t} \quad (r_{w} \leqslant r < r_{s}) \\
-\frac{k_{h}}{\gamma_{w}} \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u_{r}}{\partial r} \right) = \frac{\partial \varepsilon_{v}}{\partial t} \quad (r_{s} \leqslant r \leqslant r_{e})$$
(3)

$$2\pi r_{\rm w} d_z \frac{k_{\rm s}}{\gamma_{\rm w}} \frac{\partial u_r}{\partial r} \bigg|_{r=r_{\rm w}} = -\pi r_{\rm w}^2 d_z \frac{k_{\rm w}}{\gamma_{\rm w}} \frac{\partial^2 u_{\rm w}}{\partial z^2}$$

或

$$\frac{\partial^2 u_{\rm w}}{\partial z^2} = -\frac{2k_{\rm s}}{r_{\rm w}k_{\rm w}} \frac{\partial u_r}{\partial r} \bigg|_{r=r_{\rm w}}$$
(4a)

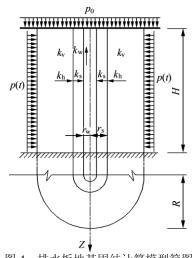


图 4 排水板地基固结计算模型简图 g.4 Calculation model diagram of drainage foundation consolidation

边界条件及初始条件为

$$u_{\mathbf{r}}\big|_{r=r_{\mathbf{b}}} = p(t) \tag{4b}$$

$$u_{\mathbf{w}}\big|_{z=0} = 0 \tag{4c}$$

$$\frac{\partial u_{\rm w}}{\partial z} = 0 \tag{4d}$$

$$\overline{u_{\rm r}}\Big|_{0} = u_0 = p_0 \tag{4e}$$

$$u_{r}\big|_{r=r_{w}}=u_{w} \tag{4f}$$

$$k_{\rm h} \left. \frac{\partial u_{\rm e}}{\partial r} \right|_{r=r_{\rm e}} = k_{\rm s} \left. \frac{\partial u_{\rm s}}{\partial r} \right|_{r=r_{\rm e}}$$
 (4g)

$$u_{s}\big|_{r=r_{s}}=u_{e}\big|_{r=r_{s}} \tag{4h}$$

式中: γ_w 为水容重(kN/m³), E_s 为土体体积压缩模量(MPa), u_r 为仅考虑径向渗流时影响区内任一点孔压(MPa), $\overline{u_r}$ 为仅考虑径向渗流时影响区内任一点平均孔压(MPa)。其中式(4e)为初始条件,其余均为边界条件。

3 求解过程

对于式(3)两边关于 r 积分,并利用边界条件(式(4b), (4f), (4g), (4h)), 可得

$$\frac{\partial u_{r}}{\partial r} = \begin{cases}
-\frac{\mathbf{r}}{2} \frac{\gamma_{w}}{k_{s}} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{p(t) - u_{w}}{\eta r} + \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{s}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\mathbf{r}}{2} \frac{\gamma_{w}}{k_{h}} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{p(t) - u_{w}}{\eta r} + \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{s}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{s}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{s}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{s}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{s}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{s}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{s}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{s}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{w}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{s}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{w}^{2} - r_{e}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{w}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{w}^{2} - r_{w}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
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-\frac{\gamma_{w}}{k_{s}} \frac{r_{w}^{2} - r_{w}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} + \frac{\gamma_{w}}{k_{h}} \frac{r_{w}^{2} - r_{w}^{2}}{4} \frac{\partial \varepsilon_{v}}{\partial t} \\
-\frac{\gamma_{$$

$$\begin{aligned} &-\frac{\gamma_{\rm w}}{k_{\rm s}} \left(\frac{r^2-r_{\rm w}^2}{4} + \frac{r_{\rm w}^2-r_{\rm s}^2}{4\eta} \ln \frac{r}{r_{\rm w}}\right) \frac{\partial \varepsilon_{\rm v}}{\partial t} + \\ &-\frac{\gamma_{\rm w}}{k_{\rm h}} \left(\frac{r_{\rm s}^2-r_{\rm e}^2}{4\eta} \ln \frac{r}{r_{\rm w}}\right) \frac{\partial \varepsilon_{\rm v}}{\partial t} + \ln \frac{r}{r_{\rm w}} \frac{p(t)-u_{\rm w}}{\eta} + u_{\rm w} \\ &-(r_{\rm w} \leqslant r < r_{\rm s}) \end{aligned}$$

$$u_r = \begin{cases} &-\frac{\gamma_{\rm w}}{k_{\rm s}} \left[\frac{r_{\rm s}^2-r_{\rm w}^2}{4} + \frac{r_{\rm w}^2-r_{\rm s}^2}{4\eta} \left(\frac{k_{\rm s}}{k_{\rm h}} \ln \frac{r}{r_{\rm s}} + \ln \frac{r_{\rm s}}{r_{\rm w}}\right)\right] \frac{\partial \varepsilon_{\rm v}}{\partial t} + \\ &-\frac{\gamma_{\rm w}}{k_{\rm h}} \left[\frac{r^2-r_{\rm s}^2}{4} + \frac{r_{\rm s}^2-r_{\rm e}^2}{4\eta} \left(\frac{k_{\rm s}}{k_{\rm h}} \ln \frac{r}{r_{\rm s}} + \ln \frac{r_{\rm s}}{r_{\rm w}}\right)\right] \frac{\partial \varepsilon_{\rm v}}{\partial t} + \\ &-\left(\frac{k_{\rm s}}{k_{\rm h}} \ln \frac{r}{r_{\rm s}} + \ln \frac{r_{\rm s}}{r_{\rm w}}\right) \frac{p(t)-u_{\rm w}}{\eta} + u_{\rm w} \quad (r_{\rm s} \leqslant r \leqslant r_{\rm e}) \end{cases}$$

其中,

$$\eta = \frac{k_{\rm s}}{k_{\rm h}} \ln n + \left(1 - \frac{k_{\rm s}}{k_{\rm h}}\right) \ln s_0$$

 $s_0 = r_{\rm s}/r_{\rm w}$, $n = r_{\rm e}/r_{\rm w}$,

式中: s_0 为涂沫比, n 为井径比。

地基中任一深度的平均孔压可表示为

$$\overline{u_r} = \frac{1}{\pi (r_e^2 - r_w^2)} \int_{r_w}^{r_e} 2\pi r u_r d_r$$
 (7)

将式(6)代入式(7),积分得

$$\overline{u_{\rm r}} = \frac{\gamma_{\rm w} r_{\rm e}^2}{4k_{\rm h}} F_{\rm a} \frac{\partial \varepsilon_{\rm v}}{\partial t} + F_{\rm b} p(t) + F_{\rm c} u_{\rm w}$$
 (8)

其中,

$$F_{a} = \frac{1}{n^{2}(n^{2} - 1)} \left[\frac{1}{2} (n^{2} - s_{0}^{2})^{2} \right] + \frac{n^{2} - s_{0}^{2}}{(n^{2} - 1)\eta}.$$

$$\left\{ \left[\ln s_{0} + \frac{k_{s}}{k_{h}} \left(\ln \frac{n}{s_{0}} - \frac{1}{2} \right) \right] + \frac{s_{0}^{2}}{2n^{2}} \left(\frac{k_{s}}{k_{h}} - 1 \right) + \frac{1}{2n^{2}} \right\} +$$

$$\frac{k_{h}}{k_{s}} \left\{ \frac{1}{n^{2}(n^{2} - 1)} \left(\frac{1}{2} s_{0}^{4} + n^{2} - s_{0}^{2} n^{2} - \frac{1}{2} \right) + \frac{s_{0}^{2} - 1}{(n^{2} - 1)\eta}.$$

$$\left[\ln s_{0} + \frac{k_{s}}{k_{h}} \left(\ln \frac{n}{s_{0}} - \frac{1}{2} \right) \right] + \frac{s_{0}^{2} - 1}{(n^{2} - 1)\eta}.$$

$$\left[\frac{s_{0}^{2}}{2n^{2}} \left(\frac{k_{s}}{k_{h}} - 1 \right) + \frac{1}{2n^{2}} \right] \right\}$$

$$F_{b} = \frac{1}{\eta(n^{2} - 1)}.$$

$$\left[\frac{1}{2} (1 - s_{0}^{2}) + n^{2} \ln s_{0} + \frac{k_{s}}{k_{h}} \left(n^{2} \ln s_{0} + \frac{1}{2} s_{0}^{2} - \frac{1}{2} n^{2} \right) \right]$$

 $F_{c} = 1 - F_{b}$

由式(5)可推导出式(4)写为

$$\frac{\partial^2 u_{\rm w}}{\partial z^2} = \frac{\gamma_{\rm w}}{k_{\rm w}} \left(1 + \frac{1 - s_0^2}{2\eta} + \frac{k_{\rm s}}{k_{\rm h}} \frac{n^2 - s_0^2}{2\eta} \right) \frac{\partial \varepsilon_{\rm v}}{\partial t} - \frac{2k_{\rm s}}{k_{\rm w} \eta r_{\rm w}^2} p(t) + \frac{2k_{\rm s}}{k_{\rm w} \eta r_{\rm w}^2} u_{\rm w} \tag{9a}$$

令

$$\alpha = \frac{2k_{\rm s}}{k_{\rm w}\eta r_{\rm w}^2} \tag{9b}$$

结合式(2), (8), (9a), (9b)可得

$$\frac{\partial \overline{u_{\rm r}}}{\partial t} = -\lambda (\overline{u_{\rm r}} - F_{\rm b} p(t) - F_{\rm c} u_{\rm w}) \tag{10}$$

$$\frac{\partial^2 u_{\rm w}}{\partial z^2} = -\rho^2 (\overline{u_{\rm r}} - F_{\rm b} p(t) - F_{\rm c} u_{\rm w}) - \alpha p(t) + \alpha u_{\rm w} \quad (11)$$

其中,

(6)

$$\lambda = \frac{4E_{s}k_{h}}{\gamma_{w}F_{a}r_{e}^{2}} = \frac{16c_{h}}{d_{e}^{2}F_{a}}$$

$$c_{h} = k_{h}E_{s}/r_{w}$$

$$\rho^{2} = \frac{4k_{h}}{k_{w}F_{a}r_{e}^{2}} \left(1 + \frac{1 - s_{0}^{2}}{2\eta} + \frac{k_{s}}{k_{h}} \frac{n^{2} - s_{0}^{2}}{2\eta}\right)$$

$$\alpha = \frac{2k_{s}}{k_{w}\eta r_{w}^{2}}$$

从式(10),(11)中消去 $\overline{u_{r}}$,得

$$\frac{\partial^{3} u_{w}}{\partial z^{2} \partial t} + \lambda \frac{\partial^{2} u_{w}}{\partial z^{2}} - (\rho^{2} F_{c} - \alpha) \frac{\mathrm{d} p(t)}{\mathrm{d} t} - (\rho^{2} F_{c} + \alpha) \frac{\partial u_{w}}{\partial t} + \lambda \alpha p(t) - \lambda \alpha u_{w} = 0$$
(12)

$$\stackrel{\text{def}}{=} 0 \leqslant t \leqslant t_1$$
 时, $\frac{\mathrm{d}p(t)}{\mathrm{d}t} = \frac{p'}{t_1}$, $p(t) = \frac{p'}{t_1}t$ 。式(12)

变为

$$\frac{\partial^{3} u_{w}}{\partial z^{2} \partial t} + \lambda \frac{\partial^{2} u_{w}}{\partial z^{2}} - (\rho^{2} F_{c} - \alpha) \frac{p'}{t_{1}} - (\rho^{2} F_{c} + \alpha) \frac{\partial u_{w}}{\partial t} + \lambda \alpha \frac{p'}{t_{1}} t - \lambda \alpha u_{w} = 0$$
(13)

欲求式(13)的解,要先求其齐次方程的解:

$$\frac{\partial^3 u_{\rm w}}{\partial z^2 \partial t} + \lambda \frac{\partial^2 u_{\rm w}}{\partial z^2} - (\rho^2 F_{\rm c} + \alpha) \frac{\partial u_{\rm w}}{\partial t} - \lambda \alpha u_{\rm w} = 0 \quad (14)$$

对 uw 进行 Laplace 积分变换:

$$\phi(z, s) = \int_0^\infty u_{\rm w}(z, t) e^{-st} dt$$
 (15)

利用 Laplace 微分性质,对式(14)进行 Laplace 积分变换得

$$(s+\lambda)\frac{\mathrm{d}^2\phi(z, s)}{\mathrm{d}z^2} - [(\rho^2 F_c + \alpha)s + \lambda\alpha]\phi(z, s) = 0$$
(16)

得到 $\phi(z, s)$ 的通解为

$$\phi(z, s) = A_1 \exp(\delta z) + A_2 \exp(-\delta z)$$
 (17)

其中,

$$\delta^2 = \frac{(\rho^2 F_c + \alpha)s + \lambda \alpha}{s + \lambda}$$

$$A_1 = \frac{p'}{1 + \exp(2\delta l)}, \quad A_2 = \frac{p' \exp(2\delta l)}{1 + \exp(2\delta l)}$$

式中: s 为 Laplace 积分变换参数。对 $\phi(z, s)$ 做 Laplace 逆变换:

$$u_{w}(z, t) = \frac{1}{2\pi j} \int_{\beta - j\infty}^{\beta + j\infty} \left[A_{l} \exp(\delta z) + A_{2} \exp(-\delta z) \right] e^{st} ds$$
(18)

则式(13)的通解可由式(18)及(13)的特解叠加得到:

$$u_w(z, t) = \frac{1}{2\pi i} \int_{\beta - j\infty}^{\beta + j\infty} [A_l \exp(\delta z) +$$

$$A_2 \exp(-\delta z) e^{st} ds + \varphi(t)$$
 (19)

其中,

$$\varphi(t) = \exp(-Dt) + Et + F$$

$$D = \frac{\lambda \alpha}{\rho^2 F_c + \alpha}, \quad E = \frac{p'}{t_1}, \quad F = \frac{-p'}{\lambda \alpha t_1} \rho^2$$

将式(18)代入式(11)中,得

$$\overline{u_r} = \left(F_c + \frac{\alpha - \delta^2}{\rho^2}\right) u_w + \left(F_b - \frac{\alpha}{\rho^2}\right) p(t) + \left(F_c + \frac{\alpha}{\rho^2}\right) \varphi(t)$$
(20)

$$A_{1} = \frac{p'}{1 + \exp(2\delta H)}, A_{2} = \frac{p' \exp(2\delta H)}{1 + \exp(2\delta H)}$$

由式(2),(6),(19)及(20),可得到满足基本方程及其求解条件的精确解:

$$u_{r} = \left[\frac{k_{h}}{k_{s}} \left(\frac{r^{2}}{r_{e}^{2}} - \frac{1}{n^{2}}\right) + \ln\frac{r}{r_{w}} \left(\frac{k_{h}}{k_{s}} \frac{1 - s_{0}^{2}}{\eta n^{2}} + \frac{s_{0}^{2} - n^{2}}{\eta n^{2}}\right)\right].$$

$$\left\{\frac{1}{\lambda F_{a}} \left(F_{c} + \frac{\alpha - \delta^{2}}{\rho^{2}}\right) s u_{w1} + \frac{F_{b} + F_{c}}{\lambda F_{a}} \frac{p'}{t_{1}} - \frac{\alpha}{\rho^{2} F_{a}} \exp\left[-\frac{\lambda \alpha}{(\rho^{2} F_{c} + \alpha)} t\right]\right\} + \left(1 - \frac{1}{\eta} \ln\frac{r}{r_{w}}\right).$$

$$\left[u_{w1} + \varphi(t)\right] + \ln\frac{r}{r_{w}} \frac{p(t)}{n} \quad (r_{w} \leqslant r \leqslant r_{s}) \tag{21}$$

$$u_{r} = \left\{ \frac{k_{h}}{k_{s}} \left[\frac{s_{0}^{2} - 1}{n^{2}} + \frac{1 - s_{0}^{2}}{\eta n^{2}} \left(\frac{k_{s}}{k_{h}} \ln \frac{r}{r_{s}} + \ln \frac{r_{s}}{r_{w}} \right) \right] + \left[\frac{r^{2}}{r_{e}^{2}} - \frac{s_{0}^{2}}{n^{2}} + \frac{s_{0}^{2} - n^{2}}{\eta n^{2}} \left(\frac{k_{s}}{k_{h}} \ln \frac{r}{r_{s}} + \ln \frac{r_{s}}{r_{w}} \right) \right] \right\}.$$

$$\left\{ \frac{1}{\lambda F_{a}} \left(F_{c} + \frac{\alpha - \delta^{2}}{\rho^{2}} \right) s u_{w1} + \frac{F_{b} + F_{c}}{\lambda F_{a}} \frac{p'}{t_{1}} - \frac{\alpha}{\lambda F_{a}} \exp \left(-\frac{\lambda \alpha}{(\rho^{2} F_{c} + \alpha)} t \right) \right\} + \left[1 - \frac{1}{\eta} \left(\frac{k_{s}}{k_{h}} \ln \frac{r}{r_{s}} + \ln \frac{r_{s}}{r_{w}} \right) \right] \left[u_{w1} + \varphi(t) \right] + \frac{p(t)}{\eta} \left(\frac{k_{s}}{k_{h}} \ln \frac{r}{r_{s}} + \ln \frac{r_{s}}{r_{w}} \right) \left(r_{s} \leqslant r \leqslant r_{e} \right)$$

$$(22)$$

进而可得地基任一深度的径向固结度 Ur为

$$U_{r} = 1 - \frac{\overline{u_{r}}}{u_{0}} = 1 - \left[\left(\frac{F_{c}\rho^{2} + \alpha - \delta^{2}}{u_{0}\rho^{2}} u_{w1} + \frac{F_{b}\rho^{2} - \alpha}{u_{0}\rho^{2}} p(t) \right) + \frac{F_{c}\rho^{2} + \alpha}{u_{0}\rho^{2}} \varphi(t) \right]$$
(23)

地基平均固结度 U_r 为

$$\overline{U_{r}} = 1 - \frac{\left(\frac{1}{H} \int_{0}^{H} \overline{u_{r}} d_{z}\right)}{u_{0}} = 1 - \left\{\frac{F_{c} \rho^{2} + \alpha - \delta^{2}}{H u_{0} \rho^{2}} \frac{1}{2\pi j}\right\}$$

$$\int_{\beta - j\infty}^{\beta + j\infty} \frac{1}{\delta} \left[A_{l} \exp(\delta H) - A_{l} - A_{2} \exp(-\delta H) + A_{2}\right] e^{st} ds + \frac{F_{b} \rho^{2} - \alpha}{u_{0} \rho^{2}} p(t) + \frac{F_{c} \rho^{2} + \alpha}{u_{0} \rho^{2}} \varphi(t)$$

$$(24)$$

式(19)~(24)即为增压为变量的等应变条件下增压式真空预压单层均质土地基径向固结问题的解答。

4 两种方法对比

分离变量法求得解析解:

$$u_{w} = \frac{2}{MF_{c}} \frac{\lambda - B_{r}}{\lambda} \left[u_{0} - \left(F_{c} + \frac{\alpha}{\rho^{2}} \right) F \right].$$

$$\sum_{m=0}^{\infty} A \sin \frac{Mz}{H} \exp(-B_{r}t) + \varphi(t) \qquad (25)$$

$$\overline{u_{r}} = \frac{2}{M} \left[u_{0} - \left(F_{c} + \frac{\alpha}{\rho^{2}} \right) F \right] \sum_{m=0}^{\infty} \sin \frac{Mz}{H} \exp(-B_{r}t) + \left(F_{b} - \frac{\alpha}{\rho^{2}} \right) p(t) + \left(F_{b} + \frac{\alpha}{\rho^{2}} \right) \varphi(t) \qquad (26)$$

$$u_{r} = \left[\frac{k_{h}}{k_{s}} \left(\frac{1}{n^{2}} - \frac{r^{2}}{r_{c}^{2}}\right) + \ln \frac{r}{r_{w}} \left(\frac{k_{h}}{k_{s}} \frac{s_{o}^{2} - 1}{\eta n^{2}} + \frac{n^{2} - s_{o}^{2}}{\eta n^{2}}\right)\right].$$

$$\left\{\frac{2}{M} \frac{B_{r}}{\lambda F_{a}} \left[u_{o} - \left(F_{c} + \frac{\alpha}{\rho^{2}}\right)F\right] \sum_{m=0}^{\infty} \sin \frac{Mz}{H} \exp(-B_{r}t) + \frac{F_{b} + F_{c}}{-\lambda F_{a}} \frac{p'}{t_{1}} + \frac{\alpha}{F_{a}(\rho^{2}F_{c} + \alpha)} \exp\left(-\frac{\lambda \alpha}{\rho^{2}F_{c} + \alpha}t\right)\right\} + \left(1 - \frac{1}{\eta} \ln \frac{r}{r_{w}}\right) \left\{\frac{2}{MF_{c}} \frac{\lambda - B_{r}}{\lambda} \left[u_{o} - \left(F_{c} + \frac{\alpha}{\rho^{2}}\right)F\right]\right\}.$$

$$\sum_{m=0}^{\infty} \sin \frac{Mz}{H} \exp(-B_{r}t) + \varphi(t) + \frac{r}{\eta n^{2}} \left(\frac{k_{s}}{k_{h}} \ln \frac{r}{r_{s}} + \ln \frac{r_{s}}{r_{w}}\right) + \frac{r_{s}}{\eta n^{2}} \left(\frac{k_{s}}{k_{h}} \ln \frac{r}{r_{s}} + \ln \frac{r_{s}}{r_{w}}\right) + \frac{r_{s}}{\eta n^{2}} \left(\frac{k_{s}}{k_{h}} \ln \frac{r}{r_{s}} + \ln \frac{r_{s}}{r_{w}}\right) + \frac{r_{s}}{\eta n^{2}} \left(\frac{k_{s}}{k_{h}} \ln \frac{r}{r_{s}} + \ln \frac{r_{s}}{r_{w}}\right) + \frac{r_{s}}{\eta n^{2}} \left(\frac{k_{s}}{k_{h}} \ln \frac{r}{r_{s}} + \ln \frac{r_{s}}{r_{w}}\right) + \frac{r_{s}}{\eta n^{2}} \left(\frac{k_{s}}{\eta n^{2}} \ln \frac{r}{\eta n^{2}} + \frac{n^{2} - s_{o}^{2}}{\eta n^{2}} \left(\frac{k_{s}}{\eta n^{2}} \ln \frac{r}{\eta n^{2}} + \frac{n^{2} - s_{o}^{2}}{\eta n^{2}} + \frac{n^{2} - s_{o}^{2}}{\eta n^{2}} \left(\frac{k_{s}}{\eta n^{2}} \ln \frac{r}{\eta n^{2}} + \frac{n^{2} - s_{o}^{2}}{\eta n^{2}} + \frac{n^{2} - s$$

地基平均固结度 ॥ 为

$$\overline{U_{r}} = 1 - \frac{\frac{1}{H} \int_{0}^{H} \overline{u_{r}} d_{z}}{u_{0}} = 1 - \left[\frac{2}{M^{2}} \left(1 - \frac{F_{c} \rho^{2} + \alpha}{u_{0} \rho^{2}} F \right) \right] \cdot (30)$$

$$\sum_{m=0}^{\infty} \exp(-B_{r}t) + \frac{F_{b} \rho^{2} - \alpha}{u_{0} \rho^{2}} p(t) + \frac{F_{c} \rho^{2} + \alpha}{u_{0} \rho^{2}} \varphi(t) \right]$$

式(25)~(30)即为利用分离变量法得到了等应

变且增压为变量条件下,增压式真空预压单层均质 土地基径向固结问题的解。

排水板等效砂井半径 $r_{\rm w}$ = 0.035 m, 涂抹区半径 $r_{\rm s}$ = 0.07 m, 涂抹比 s = 2。土层参数见表 1。

表 1 土层参数 Table 1 Soil parameters

土体区域	压缩模量 E _s /MPa	径向渗透系数 k _h /(m•s ⁻¹)	排水板渗透系数 $k_{\rm w}/({ m m} \cdot { m s}^{-1})$
影响区	1.84	2×10 ⁻⁹	1.75×10^{-5}
涂抹区	1.84	1×10^{-9}	1.75×10^{-5}

将表 1 中参数代入式(30)中,可得固结度-时间因子曲线图如图 5 所示。

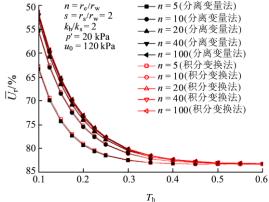


图 5 增压为变量时 2 种方法径向平均固结度 - 时间因子 曲线

Fig.5 $U_{\rm r}$ - $T_{\rm h}$ curves of the two methods when boost change over time

图 5 中黑色曲线为利用分离变量法求解得到的 径向平均固结度 - 时间因子曲线,而红色空心曲线 则为积分变换法求解得到的径向平均固结度 - 时间 因子曲线。从图 5 中可以看出当增压为变量时,两 种方法所求得的径向平均固结度随时间因子和井径 比的变化是完全一致的,数值上也相差甚小。

K. Terzaghi^[21]给出的瞬时加载条件下地基竖向固结理论解见下式:

$$\overline{U_{v}} = 1 - \sum_{v=1}^{\infty} \frac{2}{M^{2}} \exp(-M^{2}T_{v})$$
 (31)

其中,

$$M = (2m-1)\pi/2$$

式中: T_v 为竖向固结时间因子。

根据 Carrilo 定理^[22],可由竖向固结与径向固结综合起来考虑地基整体平均固结度,地基平均固结度可按下式计算:

$$\overline{U} = 1 - (1 - \overline{U}_r)(1 - \overline{U}_u) \tag{32}$$

式中: \overline{U} , 为径向平均固结度, \overline{U} , 为竖向平均固结

度。

将上述参数代入式(31),(32)中可得相应考虑竖向和径向固结的地基平均固结度,并绘制地基平均固结度-时间因子曲线如图 6 所示。

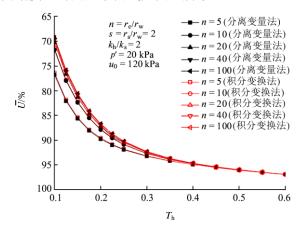


图 6 增压为变量时 2 种方法的地基平均固结度 - 时间 因子曲线

Fig. 6 \bar{U} - T_h curves of the two methods when boost change over time

图 6 中,黑色曲线为利用分离变量法求解得到的地基平均固结度 - 时间因子曲线,而红色空心曲线则为积分变换法求解得到的地基平均固结度 - 时间因子曲线。从图 6 中可以看出当增压为变量时,两种方法所求得的地基平均固结度随时间因子和井径比的变化是完全一致的,数值上也相差甚小。

由此可知,积分变换法求解增压式真空预压处理单层均质土地基固结度解析解问题是可行的,从其固结度 - 时间因子曲线可以看出,随着时间因子 T_h 的增大,固结度总体呈快速增长趋势;随着井径比 n 增大,固结速率逐渐减缓,与分离变量法所得结果几乎一致。

5 讨论

从图 5,6 结果来看,分离变量和积分变换得到的解十分吻合,这是因为在偏微分方程求解过程中,分离变量法和积分变换法有一个共同特点:都是将偏微分方程化为常微分方程,只不过转换方式不同而已。

(1) 积分变换法的思想是先找到一种函数 e^{-st} ,让函数 e^{-st} 与待求函数 u_w 积分,使得待求函数 u_w 中自变量 t 成为特征值 s。由于函数 e^{-st} 已知,所以待求导数 u_w 与函数 e^{-st} 的导数和积分关系都已知,可以代入待解偏微分方程,从而使得偏微分方程中

的变量 t 可以用特征值 s 替代,从而得到变换域内的待求函数 u_w ,再进行逆变换即可得到物理域内待求函数 u_w 的解答^[23]。

而分离变量法正好相反,分离变量的思想为: 先假设原函数 u_w 可分离变量为假设函数 Γ ,通过对偏微分方程进行分离变量,得到假设函数 Γ ,再把函数 Γ 看成广义傅里叶级数的基函数,然后通过边界条件和初始条件求解出假设函数 Γ 的系数; 对系数和基函数求和叠加,从而求得原函数 u_w 。

- (2) 从所求得函数 u_w 的形式来看,分离变量得到的是离散解,而积分变换得到的是连续解。
- (3) 当变量区间为 (-∞,+∞), 可考虑采用傅里 叶变换; 当变量区间为 (0,+∞), 可考虑采用拉普拉 斯变换; 当变量区间为有限 (a, b), 可考虑有限傅 里叶变换, 即分离变量法。

6 结 论

- (1) 随着时间因子 T_h 的增大,固结度总体呈快速增长的趋势;随着并径比n 增大,固结速率减缓。
- (2) 对比积分变换法和分离变量法所得解析解,两种方法计算所得结果一致,数值几乎无差别。
- (3) 积分变换法与分离变量法求解原理有所不同; 积分变换法是先将偏微分方程转换为更易求解的函数形式, 在变换域内求解之后进行逆变换得到原函数的解答; 而分离变量法是求函数的特解, 将特解叠加得到原函数的通解; 2 种方法殊途同归, 求得解析解数值相差甚小。

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